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# MATRIX ISOLATION OF BORON AND CARBON VAPOR. CONTROL OF CLUSTER FORMATION DURING PREPARATION AND ANNEALING

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#### **Abstract**

The formation and recombination of small boron/carbon clusters were studied by quantitative Fourier transform infrared (FTIR) spectroscopy. Samples were prepared by evaporation of powder mixtures from a resistively heated oven (3000 K) and isolated at 10 K in solid argon.  $C_3$  and  $BC_2$  disappeared entirely over the course of several annealings. Yields of all the larger clusters ( $B_1C_{n-1}$ ,  $n \ge 4$ , J = 0, 1, 2) increased during annealing. Clusters with  $n \ge 4$  are linear, with the boron atoms of the J = 1 and 2 clusters capping the ends of carbon chains;  $B_2C_{n-2}$  clusters are inert. During annealing,  $C_n$  and  $BC_{n-1}$  clusters grow by adding mobile atom, dimer, or trimer to the carbon end of the chain. Dramatic growth of  $B_2C_2$  during the first annealing is consistent with the presence of B atoms and/or  $C_1$  BC molecule in the initial matrices.

#### Introduction

Under the auspices of the U. S. Air Force High Energy Density Matter (HEDM) Program [1], we have pursued the goal of production of a cryogenic matrix with an atom density of  $\sim 1 \times 10^{21}$  atoms cm<sup>-3</sup>, which amounts to  $\sim 5$  mole percent atoms in an inert solid, or an equivalent atom gas pressure of  $\sim 60$  atm at 298 K. This and our companion paper [2] describe qualitative and quantitative analysis of  $B_1C_{n-1}$  clusters in prototypical HEDM. Previously, the first conclusive identifications of small boron clusters, BC<sub>2</sub>, B<sub>2</sub>C, BC<sub>3</sub>, and  $B_2C_2$  were reported [3-5]. Observed FTIR fundamentals and isotope shifts were matched with predictions of *ab-initio* theory in those studies.

Verhaegen, Stafford, and Drowart [6] published the first experimental evidence of B<sub>1</sub>C<sub>n-7</sub> 35-years ago. Mass spectrometry was used to measure boron and boron carbide vaporization from graphite Knudsen cells heated by electron bombardment to ~ 2500 K. The vapor pressure of boron at 2200 K was ~ 0.008 torr, 94% atoms, 5% BC<sub>2</sub>, 1% B<sub>2</sub>C, and 0.07% BC. At the same temperature, boron carbide had about one-half the vapor pressure with 96% atoms, 4% BC<sub>2</sub>, 0.1% B<sub>2</sub>C. Distributions of anions and cations of B<sub>x</sub>C<sub>y</sub> with up to 17-atoms have been measured in plumes of laser vaporized boron carbide targets by Becker and Dietze [7]. Their mass spectrometric analysis showed that B<sub>2</sub>C<sup>-</sup>, B<sub>2</sub>C<sup>+</sup>, BC<sub>2</sub><sup>-</sup>, and BC<sub>2</sub><sup>+</sup> were major components of the plume. Matrix isolated BC<sub>2</sub> has been produced by oven and laser vaporization of carbon/boron mixtures and targets and measured by FTIR [3,8] and ESR [9,10] spectroscopy. Most recently, Wyss, Grutter and Maier [11] isolated BC<sub>2</sub><sup>-</sup> and BC<sup>-</sup> in neon matrices by sputtering a boron carbide target with Cs<sup>+</sup> and trapping the mass-selected anions in neon matrices. Neutrals were produced by electron photodetachment with uv radiation. Electronic transitions of BC<sub>2</sub> and BC and their anions were studied. The FTIR

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spectrum of the previously identified [3]  $v_2(a_1)$  fundamental of  $BC_2$  was also observed. The infrared spectrum of BC has never been measured, but a CCSD(T)/cc-pVTZ calculation with quartic force field predicted [12] its vibrational fundamental at 1148 cm<sup>-1</sup> with intensity of 43 km mol<sup>-1</sup>. Of the pure boron clusters, only  $B_3$  has been experimentally observed. Li, Van Zee and Weltner [13] identified all four isotopomers of  $B_3$  in argon matrices by FTIR spectroscopy and found that experimental frequencies matched *ab-initio* frequencies of its cyclic geometry. Although *ab-initio* studies of  $B_n$ ,  $n \ge 4$ , have been carried out, only geometries and energies have been reported [14,15]. Unlike  $C_n$ ,  $B_n$  for  $n \ge 5$  are predicted to have stable 3-dimensional geometries.

Carbon clusters were recently reviewed by Van Orden and Saykally [16]. Cyclic  $C_6$  [17, 18], cyclic  $C_8$  [19, 20], linear  $C_{11}$  [21], have recently been identified by their FTIR isotopomer spectroscopy. Most recently, electronic spectra of linear  $C_{17}$ ,  $C_{19}$ , and  $C_{21}$  [22] and cyclic  $C_{10}$  and cyclic  $C_{12}$  [23] isolated in neon matrices were reported, and Raman spectra of linear  $C_{16}$ ,  $C_{18}$ , and  $C_{20}$  were reported [24].

Figure 1 shows a Pascal's triangle representation of the boron/carbon system. In the limit of statistical recombination of well mixed atoms, the relative yields of  $B_1C_{n-1}$  clusters of given n is given by

$$\rho(B_{J}C_{n-J})/\rho(C_{n}) = n!/J!(n-J)/[B/C]^{J},$$
(1)

where [B/C] is the molar ratio of boron to carbon, and  $\rho(i)$  is the column density of species i.

# Experimental

Mixtures of boron and carbon powders were evaporated with a resistively heated oven to temperatures ~ 3000 K, and vapors were co-condensed with ~ 1000:1 excess argon at  $10^{\circ}$  K to form ~ 60  $\mu$  thick matrices in about one-hour. Beer's Law enabled quantitative analysis of the various species,  $\rho(i) = 2.303 \ A_i(experiment)/I_i(theory)$ , where  $\rho(i)$  is the average column density, which is the average value of the product of the species density and matrix thickness.  $A_i(experiment)$  is the integrated intensity of the infrared absorbance of species i, equal to  $-log_{10} \tau_i(v)$ , where  $\tau_i(v)$  is the transmittance at absorption frequency v.  $I_i(theory)$  is the infrared absorption intensity of species i, which by necessity must be obtained by ab-initio calculation because experimental values are not available at this time. Fig. 1 indicates values of  $I_i(theory)$  calculated by DFT/cc-pVDZ [25-28].

# **Results and Discussion**

Annealing of pure carbon matrices led to disappearance of atoms, dimers and trimers from the initial deposits. The identified products of recombination included cyclic  $C_6$  and cyclic  $C_8$ , and linear  $C_n$  molecules,  $4 \le n \le 13$ . After annealing to constant composition, the total measurable carbon was found to be as much as twice the initial measurable carbon, indicating that about 50 % of the initially deposited carbon was isolated as atoms and dimers. The cyclic molecules, the dominant condensation products, amounted to as much as  $\sim 60\%$  of total carbon in the fully annealed matrices.

In matrices containing boron, linear  $C_3$  and cyclic  $BC_2$ ,  $B_2C$ , and  $B_3$ , constituted about 80% of the total observable boron and carbon in the initially deposited matrix. The observation of approximately statistical distributions [Equation (1)] in clusters with n=3, 4, and 5 in the initially deposited matrices implies that these smaller clusters form by random condensation of well-mixed atoms, uninfluenced by their relative energies, the energies of their precursors, or preferential kinetics pathways that could otherwise distort the statistics.

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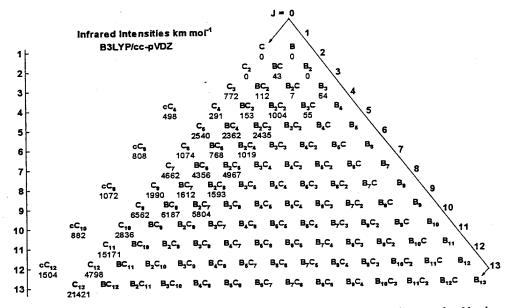


Figure 1. Representation of the  $B_1C_{n,J}$  (n=1-13, J=0-n) system by Pascal's triangle. Numbers below molecular formulas denote the <u>ab-initio</u> intensities of the most intense infrared stretching fundamentals, Refs. [25] – [28]. Cyclic carbon clusters are removed from the triangle to the left.

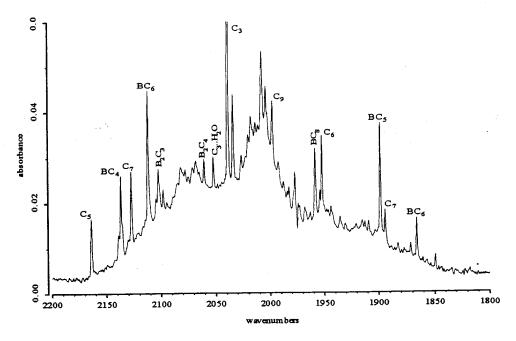


Figure 2. Survey spectrum of matrix containing carbon and boron at natural abundance with  $[B/C] \sim 1/3$  after three annealings. All of the peaks indicated grow upon annealing except  $C_3$ . Fundamentals of  $BC_{n-1}$  for n=5, 6, 7, and 9 are similarly red-shifted from fundamentals of linear  $C_n$  and their experimental absorbances are all slightly greater in this matrix. Two fundamentals of  $BC_6$  are observed at 2112 and 1866 cm<sup>-1</sup>, red-shifted from the two fundamentals of linear  $C_7$ .

à

 $B_2C_{n-2}$  clusters, once formed, were inert to further condensation upon annealing. Thus, energy loss by recombination during preparation of high energy density matrices (HEDM) is minimized in matrices with large B/C ratios. For example, when B/C ~2, any recombination that does occur during co-deposition would tend to produce  $B_2C$ , a high-energy molecule that is inert in the cryogenic environment. In pure boron HEDM, the cyclic  $B_3$  molecule may be similarly inert, which would suggest that energy loss by recombination will also stop at a  $B_3$  "island of stability." Continuing experimental work to produce five-mole percent pure boron HEDM is concentrating on development of a resistively heated high-flux boron atom source, the "boron cannon", and production of higher density matrices.

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# MATRIX ISOLATION OF BORON AND CARBON VAPOR.

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## **Abstract**

This report supplements the companion paper published in this volume [1]. Details about the FTIR quantitative analysis of the disappearance and formation of clusters during annealing of solid argon matrices containing boron and carbon species are presented.

## Introduction

The companion paper [1] summarized the literature of small boron clusters, about  $B_1C_{n-J}$  (n = 2 to 10, J = 0, 1, 2, 3) and described the formation of clusters during preparation and annealing of solid argon. This report presents additional details.

# **Experimental**

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A precision matched pair of matrices was prepared as described in Ref. [1] with two different boron isotopic compositions and natural abundance carbon. Matrix (a) contained natural abundance boron,  $^{11}B/^{10}B = 4.0$ , and matrix (b) contained boron with an inverted isotope ratio,  $^{11}B/^{10}B = 0.37$ . Each matrix had a [C/B] molar ratio of  $\sim 3.0$ , a thickness of  $\sim 60$  m, and very similar absolute amounts of each element. Each matrix was annealed repeatedly according to the same annealing protocol given in the caption of Figure 1.

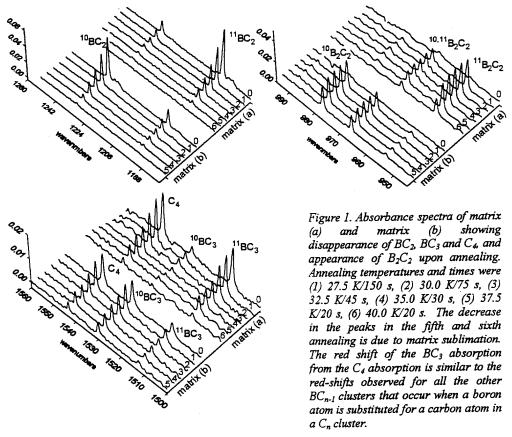
# **Results and Discussion**

Figure 1 shows the annealing behaviors of the FTIR bands belonging to  $BC_2$ ,  $BC_3$  and  $B_2C_2$  for each of the matrices. Cyclic  $BC_3$  disappears entirely when one of its carbon atoms is attacked by a radical species. One of the B-C bonds is broken, and the  $BC_3$  ring opens. The dramatic appearance of  $B_2C_2$  during the first annealing is consistent a mechanism involving recombination of BC, whose density lies below the limit of detection, or condensation of B atoms with  $BC_2$ .

Figure 2 shows the annealing behaviors of the  $^{11}B_J^{12}C^{n-1}$  clusters, n=3 to 9 and J=0,1,2 that were measured in matrix (a). Loss of clusters during the sixth annealing by matrix sublimation is generally greatest for  $C_n$  and smallest for  $B_2C_{n-2}$ , which is consistent with a stratified matrix with a higher [C/B] ratio near the subliming surface. Apparently, boron evaporates more rapidly than carbon during the early stages of oven evaporation, resulting in boron-rich layers near the substrate.

Column densities of the trimer clusters in the initially deposited matrices are in the ratio  $\rho(C_3):\rho(BC_2):\rho(B_2C):\rho(B_3)\approx 1.0:1.5:0.5:<0.05$ , where only the upper limit for  $\rho(B_3)$  may be determined because it was **not** observed. This is consistent with formation of the trimers by statistical condensation from a matrix with  $[B/C] \sim 1/3$ , which would produce a 1.0:1.0:0.33:0.05 distribution,  $\rho(B_3C_{n-1})/\rho(C_n) = n!/J!(n-J)![B/C]^J$ , Eq. (1), Ref. [1].

Agreement between distributions implies trimers form by random condensation of well-mixed atoms, uninfluenced by the relative energies of the trimers, the energies of their precursors, or preferential kinetics pathways that could otherwise distort the statistics. Linear  $C_3$  and cyclic  $BC_2$ , disappeared entirely when the matrices were repeatedly annealed to



temperatures between 25 K and 35 K, but cyclic  $B_2C$  was inert. Linear  $C_4$  and  $BC_3$  (BCCC) disappeared more slowly, and linear  $B_2C_2$  (BCCB) grew to  $\sim 95\%$  of its final value during the first annealing. Once formed,  $B_2C_2$ , like  $B_2C$ , was also inert to further reaction.

Figure 3 shows the absolute column densities of the analyzed clusters that was obtained by transformation of the Fig. 2 data with Beer's law and theoretical intensities shown in Fig. 1, Ref [1]. The BC molecule was not observed, but its upper limit column density was established to be somewhat larger than that of  $B_2C_2$ .

The sources of  $B_2C_2$  are from condensation of atom plus trimer (B + BC<sub>2</sub> but not C + B<sub>2</sub>C) or dimer + dimer (BC + BC but not B<sub>2</sub> + C<sub>2</sub>). Although BC was not observed, the upper limit of  $\rho(BC)$  is larger than  $\rho(B_2C_2)$  so that BC cannot be ruled out as a source of B<sub>2</sub>C<sub>2</sub>. The growth of B<sub>2</sub>C<sub>2</sub> is consistent with of the presence of BC and/or B in the originally deposited matrix in an amount at least as great as the growth of B<sub>2</sub>C<sub>2</sub>.

Figures 1 and 2 show that  $C_4$  and  $BC_3$  do not change during the first annealing but begin to disappear in subsequent annealings. This indicates that a source for these tetramers, atom + trimer and/or dimer + dimer, operates during the first annealing.

Figure 2 shows that linear  $C_5$ ,  $BC_4$  (BCCCC) and  $B_2C_3$  (BCCCB) and larger linear clusters all grew upon annealing. The sources of  $B_2C_3$  are dimer + trimer (BC + BC<sub>2</sub> but not  $B_2 + C_3$ ) and atom + tetramer (B + BC<sub>3</sub> but not C + B<sub>2</sub>C<sub>2</sub>). Since  $\rho(BC_2) \sim 5\rho(BC_3)$  in the initially deposited matrix, the BC + BC<sub>2</sub> source is dominant. Growth of B<sub>2</sub>C<sub>3</sub> is consistent with the presence of BC in the initially deposited matrix in an amount at least as great as the amount by which B<sub>2</sub>C<sub>3</sub> grows.

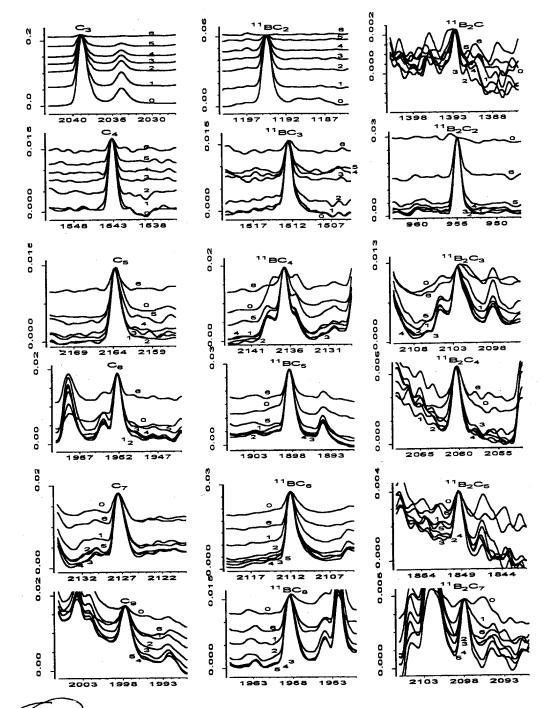


Figure 1) Annealing behaviors of  $^{11}B_J^{12}C_{nJ}$  species in matrix (a), which contained natural abundance carbon and boron,  $^{11}B_J^{10}B=80/20$ ,  $^{12}C_J^{13}C=99/1$ , with  $[C/B]\sim 3$ . Spectra labeled '0' were obtained from the originally deposited matrix, and spectra labeled '1' to '6' were obtained after successive annealings. During the 6th annealing (at 40 K for 20 seconds), almost half the matrix evaporates.

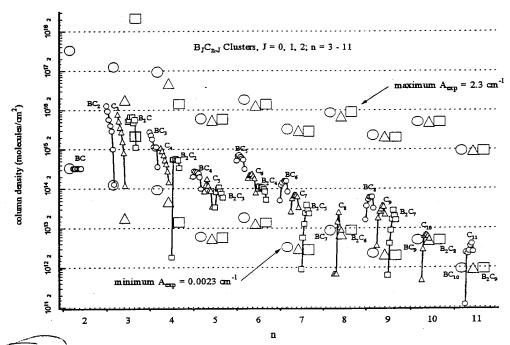


Figure 2. Distributions of  $B_1C_{n-1}$  clusters, J=0, I, 2; n=3-11. Circles, triangles and squares represent  $BC_{n-1}$ ,  $C_n$  and  $BC_{n-2}$  clusters, respectively. Large symbols denote upper and lower limits of measurement, based on a nominal minimum measurable absorbance of 0.0023 cm<sup>-1</sup>, and a maximum absorbance for linearity of Beer's law of 2.3 cm<sup>-1</sup> (1% transmittance). Small symbols denote measured quantities in the initial matrix, and in six annealed matrices. Some of the larger clusters (n=8,10,11) have not been identified,  $BC_7$ ,  $B_2C_6$ ,  $BC_9$ ,  $B_2C_8$ ,  $BC_{10}$ ,  $B_2C_9$ .

Growth of BC<sub>4</sub> occurs primarily by BC + C<sub>3</sub> rather than B + C<sub>4</sub> or C + BC<sub>3</sub> because  $\rho(C_3) \sim 10 \rho(C_4)$  and  $\rho(C_3) \sim 2 \rho(BC_3)$ . Growth of C<sub>5</sub> occurs by C + C<sub>4</sub> and C<sub>2</sub> + C<sub>3</sub>, which establishes the presence of C and/or C<sub>2</sub> in the original matrix in an amount at least as great as C<sub>5</sub> growth.

# **Conclusions**

Disappearance of triangular BC<sub>2</sub> requires breaking of one of its B-C bonds when one of its carbon atoms is attacked. The major reorganization of electronic energy involved in opening the ring appears to occur with little (< ~3 kcal mol<sup>-1</sup>) or no energy barrier, which makes this small molecule a candidate for an interesting *ab-initio* study of unusual reactivity at low temperature. Annealing kinetics of disappearance of C<sub>3</sub> and BC<sub>2</sub>, and of appearance of B<sub>2</sub>C, C<sub>4</sub>, BC<sub>3</sub>, B<sub>2</sub>O<sub>2</sub>, C<sub>5</sub>, BC<sub>4</sub>, and B<sub>2</sub>C<sub>3</sub> unequivocally establishes the presence of atoms and dimers in the originally deposited matrix. About 80% or more of the initially deposited HEDM existed as atoms, dimers and trimers. Molecules with two boron atoms are immune from radical attack and condensation during annealing.

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